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## LETTER TO THE EDITOR

# Superior mechanical properties of Nb<sub>2</sub>AsC to those of other layered ternary carbides: a first-principles study

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#### Abstract

Nb<sub>2</sub>AsC showed superior mechanical properties to those of other layered ternary carbides (Kumar *et al* 2005 *Appl. Phys. Lett.* **86** 111904). In the present density functional calculation, the underlying mechanism is interpreted by astonishing bonding features of Nb<sub>2</sub>AsC. The Nb d–As  $(p_x + p_y)$  and Nb d–As  $p_z$  bonding states locate in the same energy range as those of Nb d–C p bonding, which indicates that the Nb–As bond has similar bonding strength as the Nb–C bond does; and thereafter, Nb<sub>2</sub>AsC has improved mechanical properties compared to the others. The present reported bonding features are interestingly different from those experienced by T<sub>2</sub>AlC (T = Ti, V, Cr, and Nb), wherein the weak T–Al coupling was separated from T–C bonding states in the higher energy level by a pseudogap. This work proposes an effective method to strengthen the relative weaker T–A (A is the A-group element) bonding in layered ternary carbides.

(Some figures in this article are in colour only in the electronic version)

Layered ternary carbides with the common formula  $T_{n+1}AX_n$  (where T is a transition metal, A is an A-group element, and X is C) are promising damage-tolerant ceramics with interesting properties [1]. They exhibit merits of both metals and high-performance ceramics such as good thermal and electrical conductivity, excellent thermal shock resistance and high-temperature oxidation resistance, and microscale ductility at room temperature. Recently, optimization of mechanical properties was the most concentrated subject for TAX carbides by properly tuning the valence electron concentration (VEC) in material [2–4]. Barsoum *et al* succeeded in strengthening Ti<sub>2</sub>AlC by synthesizing the Ti<sub>2</sub>AlC<sub>0.5</sub>N<sub>0.5</sub> solid solution [2]. Later, theoretical calculations by Wang *et al* [3] and Sun *et al* [4] suggested an alternative way by tailoring the transition-metal atom in T<sub>2</sub>AlC (T = Ti, V, Cr and Nb). Increased bulk modulus and elastic

coefficient  $c_{44}$  were achieved therein. The mechanism was attributed to strengthening of T d– Al p bonding when the VEC increased [3].

The crystal structure of layered ternary carbides can be described as TC slabs in an NaCltype structure being interleaved and mirrored by close-packed A-element atomic planes. The mechanical properties of TAX carbides are dominated by the weak coupling between transitionmetal carbide/nitride slabs and the interleaved A-element atomic planes [5]. It is necessary to find an effective method to strengthen the T-A bonding in TAX material. Recently, a room-temperature synchrotron x-ray diffraction experiment was conducted on Nb<sub>2</sub>AsC in the pressure range up to 41 GPa [6]. It was reported that Nb<sub>2</sub>AsC yielded the highest bulk modulus (224 GPa) in the TAX carbides measured to date. In addition, the compressibility along the basal plane is higher than along the *c*-direction, which is different to other TAX carbides such as Ti<sub>3</sub>SiC<sub>2</sub> [7] and Ti<sub>2</sub>AlC [8]. The underlying mechanism of superior mechanical properties of Nb2AsC should be investigated, and the results of continued research in this area may play a key role in better understanding the mechanical properties of ternary TAX carbides. In the present density functional calculations, the elastic stiffness and ideal shear strength were calculated and interpreted by interesting bonding characteristics. For comparison, electronic structure analysis of Nb<sub>2</sub>AlC was also performed. Specially, we suggest that mechanical properties of TAX carbides could be more effectively enhanced by tailoring the VEC of the A-element atoms

The density functional calculations were performed using the CASTEP total energy code [9]; the Vanderbilt-type ultrasoft pseudopotential [10] and generalized gradient approximation (GGA-PW91) [11] were employed. The plane-wave basis set cutoff was set to 450 eV for all calculations. The special points sampling integration over the Brillouin zone was employed by using the Monkhorst–Pack method with a  $10 \times 10 \times 2$  special *k*-point mesh [12]. The Broyden–Fletcher–Goldfarb–Shanno (BFGS) minimization scheme [13] was used in geometry optimization. Lattice parameters, including lattice constants and internal atomic coordinates, were modified independently to minimize the total energy and interatomic forces. The tolerances for geometry optimization are: difference on total energy within  $5 \times 10^{-6}$  eV/atom, maximum ionic Hellmann–Feynman force within 0.01 eV Å<sup>-1</sup>, maximum ionic displacement within  $5 \times 10^{-4}$  Å, and maximum stress within 0.02 GPa.

To study interatomic bonding characteristics, the projected densities of states (PDOSs) were obtained within the full-potential linearized augmented plane-wave (FPLAPW) + local orbitals (lo) method implemented in the Wien2K code [14]. Separation radii between valence states and core states were 2.21, 2.50, and 1.96 au for the Nb, As, and C atoms, respectively, in Nb<sub>2</sub>AsC; but were 2.17, 2.50, and 1.92 au for the Nb, Al, and C atoms, respectively, in Nb<sub>2</sub>AlC. A  $10 \times 10 \times 2$  special *k*-point mesh was used. The total energy was converged within  $6.8 \times 10^{-6}$  eV. The spheres separate the space between core states treated as atomic orbitals and valence states treated as plane waves in the interstitial space.

The elastic stiffness determines the response of a material to an externally applied strain or stress and provides information about the bonding characteristics and structural stability. The bulk modulus measures the resistance of a material to a volume change, and the shear modulus and shear strength determine how readily a dislocation is generated and is able to move throughout a solid in response to a shear strain or stress. Furthermore, shear deformation resistance is correlated to the hardness of a material and may serve as a suitable hardness predictor [15–17]. Therefore, mechanical parameters, including bulk modulus B, shear modulus  $c_{44}$ , and ideal shear strength  $\sigma$ , were investigated in the current work.

Theoretical elastic coefficients were determined by applying a set of given homogeneous deformations with a finite value and calculating the produced stress with respect to optimizing the internal atomic coordinates [3]. The criteria for convergences of geometry optimization

**Table 1.** Computed bulk modulus *B* (in GPa), elastic constant  $c_{44}$  (in GPa) and ideal shear strength  $\tau$  (in GPa) of Ti<sub>2</sub>AlC, Nb<sub>2</sub>AlC, Cr<sub>2</sub>AlC and Nb<sub>2</sub>AsC.

Compound	В	C44	$c_{44}/B$	τ
Ti <sub>2</sub> AlC	137	111	0.810	11
Nb <sub>2</sub> AlC	183	150	0.820	14
Cr <sub>2</sub> AlC	193	147	0.762	11
Nb <sub>2</sub> AsC	219	182	0.831	19

were set as: differences in the total energy within  $1 \times 10^{-6}$  eV/atom, ionic Hellmann–Feynman forces within 0.002 eV Å<sup>-1</sup>, and the maximum ionic displacement within  $1 \times 10^{-4}$  Å.

In order to obtain the theoretical stress–strain relations, we applied a series of incremental strains to the crystal up to strains far beyond the linearity. At each step, the lattice parameters and internal atomic freedoms were simultaneously relaxed until all components of the Hellmann–Feynman stress tensor orthogonal to the applied strain converged within 0.2 GPa. The ideal shear strength was calculated as the first maximum shear stress in the theoretical stress–strain curve [17]. In the current investigation, the lattices were sheared along the  $[1\bar{2}10](0001)$  deformation path. For comparison, the elastic stiffness and ideal shear stress–shear strain relationships were calculated for the Nb<sub>2</sub>AsC, Ti<sub>2</sub>AlC, Cr<sub>2</sub>AlC and Nb<sub>2</sub>AlC compounds. The present authors have investigated the theoretical stress–strain relationships for ternary compounds Ti<sub>2</sub>AlC and Ti<sub>2</sub>AlN, and binary counterparts TiC and TiN [18]. In addition, the obtained stress–strain relationships of TiC and TiN agreed well with those reported by Jhi and collaborators [19]. The present calculation is reliable for calculating the theoretical stress–strain relationships of TiC and TiN agreed well with those reported by Jhi and collaborators [19]. The present calculation is reliable for calculating the theoretical stress–strain relationships of TiC and TiN agreed well with those reported by Jhi and collaborators [19]. The present calculation is reliable for calculating the theoretical stress–strain relationships of layered ternary compounds and their binary counterparts.

Knowledge of elastic coefficients enables us to evaluate the ratio between linear compressibility coefficients in the directions along the basal plane and the *c*-direction using [20]

$$k_c/k_a = (c_{11} + c_{12} - 2c_{13})/(c_{33} - c_{13})$$
<sup>(1)</sup>

where  $c_{ij}$  is second-order elastic coefficient. The values of  $k_c/k_a$  for Cr<sub>2</sub>AlC and Nb<sub>2</sub>AsC are 0.878 and 0.673, respectively, while the ratio is 1.139 for  $Ti_2AIC$  and 1.140 for Nb<sub>2</sub>AIC. For Cr<sub>2</sub>AlC and Nb<sub>2</sub>AsC, the stiffness along the *c*-direction is higher than that along the basal plane. In contrast, Ti2AlC and Nb2AlC experience the reverse tendency, for which the cdirection is softer. The calculated compressibility agrees well with experimental results [6]. The underlying reason for the diverse compressibility tendency could be addressed to different coupling strengths between the transition-metal carbide slabs and A-element atomic planes. In Cr<sub>2</sub>AlC and Nb<sub>2</sub>AsC, the T-A bonding strengths are suggested as being enhanced over that in  $Ti_2AIC$  and  $Nb_2AIC$ . For compounds with hexagonal symmetry, the elastic coefficient  $c_{44}$  determines the elastic resistance against an applied [1210](0001) shear strain. It has been reported that higher elastic coefficient  $c_{44}$  and, thus, improved mechanical properties (i.e., hardness) can be achieved by tuning the VEC of ternary  $T_2AIC$  (T = Ti, V, and Cr) carbides [3, 21]. The computed elastic coefficient  $c_{44}$  of Nb<sub>2</sub>AsC is listed in table 1, together with those of Ti<sub>2</sub>AlC, Nb<sub>2</sub>AlC and Cr<sub>2</sub>AlC for comparison. We note that Nb<sub>2</sub>AsC has the highest c<sub>44</sub> in these ternary compounds. This result suggests that Nb<sub>2</sub>AsC may show improved mechanical properties dominated by the shear modulus  $c_{44}$ .

Ideal shear strength is the maximum shear stress that chemical bonding can sustain under shear deformation and is of particular importance for understanding the structural stability of a solid. Theoretically, we have shown that  $Ti_2AIC$  and  $Ti_2AIN$  experience small ideal shear strengths, which leads to damage tolerance and low hardness [18]. Therein, the authors



Figure 1. Theoretical stress–strain curves of Ti<sub>2</sub>AlC, Nb<sub>2</sub>AlC, Cr<sub>2</sub>AlC and Nb<sub>2</sub>AsC under applied [1210](0001) shear strains.



**Figure 2.** Valence electron density of a slice of the  $(11\overline{2}0)$  plane in the Nb<sub>2</sub>AsC unit cell under shear deformations of (a)  $\varepsilon = 0$ , (b)  $\varepsilon = 0.1$ , (c)  $\varepsilon = 0.15$  and (d)  $\varepsilon = 0.2$ .

proposed that the ideal shear strengths of ternary compounds are much smaller than those of the binary counterparts, like TiC and TiN. Exploring enhanced ideal shear strength is expected to be important in predicting optimized mechanical properties of TAX carbides. In figure 1, we present the first-principles stress–strain curves of Ti<sub>2</sub>AlC, Nb<sub>2</sub>AlC, Cr<sub>2</sub>AlC and Nb<sub>2</sub>AsC under  $[1\bar{2}10](0001)$  shear strains. As shown in figure 1, the stress of Nb<sub>2</sub>AsC exhibits larger value than other carbides throughout the studied shear strains. The obtained ideal shear strength of Nb<sub>2</sub>AsC is about 1.7, 1.4, and 1.7 times larger than that of Ti<sub>2</sub>AlC, Nb<sub>2</sub>AlC, and Cr<sub>2</sub>AlC, respectively.

Figure 2 shows slices of valence electron density distribution in the  $(11\bar{2}0)$  atomic plane of Nb<sub>2</sub>AsC under  $[1\bar{2}10](0001)$  shear strains. When Nb<sub>2</sub>AsC is shear deformed, As atoms slip smoothly along the basal plane, and only half of the Nb–As bonds stretch and soften gradually according to the applied strain mode. In contrast, the strong Nb–C covalent bonds remain stable. The shear deformation proceeds smoothly by bond-softening rather than by an abrupt bond-breaking of the Nb–As bond. At a shear strain of 0.2, the stretched Nb–As bond reaches



Figure 3. Projected electronic density of states of (a) Nb<sub>2</sub>AlC and (b) Nb<sub>2</sub>AsC.

its limit of stability and breaks eventually. This leads to the gradual stress relaxation at large strains.

As summarized in table 1, Nb<sub>2</sub>AsC yields the highest bulk modulus B, shear modulus  $c_{44}$ , and ideal shear strength compared to Ti<sub>2</sub>AlC, Nb<sub>2</sub>AlC, and Cr<sub>2</sub>AlC. Therefore, Nb<sub>2</sub>AsC is expected to have significantly strengthened mechanical properties. It should be noted that Cr<sub>2</sub>AlC and Nb<sub>2</sub>AsC have the same average VEC in the unit cell. However, the theoretical mechanical parameters of Nb<sub>2</sub>AsC are obviously higher than those of Cr<sub>2</sub>AlC. Therefore, tailoring the A-element atoms may be more effective than tuning the VEC of transition-metal atoms to strengthen the TAX carbides. To better understand the origin of the superior mechanical properties of Nb<sub>2</sub>AsC, interatomic bonding is characterized and compared with that of Nb<sub>2</sub>AlC, a ternary carbide that differs from Nb<sub>2</sub>AsC only in element A.

The projected electronic density of states (PDOS) of Nb<sub>2</sub>AsC is illustrated in figure 3, together with those of Nb<sub>2</sub>AlC. In figure 3(a), a pseudogap separates the strong covalent Nb–C bonding states from those of Nb–Al bonding orbitals. This bonding feature is mostly observed in Ti<sub>2</sub>AlC, V<sub>2</sub>AlC, and Cr<sub>2</sub>AlC [3, 4]. However, as seen in figure 3(b), the bonding states in Nb<sub>2</sub>AsC exhibit different features: the strong Nb d–C p bonding states approximately locate between -7.0 and -3.0 eV below the Fermi level; and in addition, the Nb d–As p<sub>z</sub> and Nb d–As p<sub>x</sub> + p<sub>y</sub> bonding states, ranging from -7.0 to -2.0 eV, locate approximately in the same energy level as that of the Nb–C bond; at the same time, relatively weaker Nb d–As (p<sub>x</sub> + p<sub>y</sub>) bonding states locate in higher energy level, extending from -2.0 to -1.0 eV. The Nb d–Nb d metallic bonding dominates states near the Fermi level. Of the most interest, the bonding strengths of Nb–C and Nb–As bonds are expected to have the same scale, suggested by the



**Figure 4.** Valence electron density from different energy ranges of the DOS for Nb<sub>2</sub>AlC and Nb<sub>2</sub>AsC. The plots are in the (11 $\overline{2}0$ ) plane and the units are electron Å<sup>-3</sup>. (a) Covers energies between -7.0 and -3.0 eV of Nb<sub>2</sub>AlC, while (b) includes energies from -3.0 eV to the Fermi level. (c) and (d) are similar to (a) and (b) except for Nb<sub>2</sub>AsC.

fact of their extending in the same energy level. This shows that the bonding characteristics of Nb<sub>2</sub>AsC differ from those of Ti<sub>2</sub>AlC, V<sub>2</sub>AlC, Nb<sub>2</sub>AlC and Cr<sub>2</sub>AlC significantly. In those Al-containing ternary carbides, the bonding states associated to T–C and T–Al bonds extend in different energy levels and are separated by a noticeable pseudogap [3, 4]. Therein, the T–Al bond was definitely weaker than the T–C bond in T<sub>2</sub>AlC (T = Ti, V, Cr, and Nb).

Since the bonding characters at different energy ranges appear to be responsible for the mechanical features of the studied compounds, it is of interest to examine the charge density decomposed into different covalent bonding peaks in the DOS figures, as shown in figure 4. The plots use electron  $Å^{-3}$  as the unit of charge density. In figure 4(a), the charge density comes from states in the energy range between -7.0 and -3.0 eV for Nb<sub>2</sub>AlC. It shows clear evidence of predominant p-like orbitals on the C sites, and t<sub>2g</sub>-like orbitals on the Nb sites. Figure 4(b) illustrates the charge density originated from states in a higher energy level extending from -3.0 eV to the Fermi level for Nb<sub>2</sub>AlC. It indicates that the charge density of Al atoms has  $p_z$  symmetry, as well as the  $(e_g + t_{2g})$ -like orbitals on the Nb sites. From figures 4(a) and (b), we found that the covalent bonding in Nb<sub>2</sub>AlC is basically similar to that in  $Ti_2AlC$ ,  $V_2AlC$  and  $Cr_2AlC$ . Figure 4(c) shows the charge density corresponding to states in Nb<sub>2</sub>AsC extending in the same energy level as in figure 4(a). It is interesting to observe that charge density originated from Nb-As bonding is presented, besides those corresponding to strong Nb–C covalent bonding. The charge density in figure 4(d) derives from the states in energy level from -3.0 eV to the Fermi level of Nb<sub>2</sub>AsC. It corresponds to the double peaks in the upper energy level in figure 3(b). In contrast to the bonding features in Nb<sub>2</sub>AlC, wherein Nb-Al bonding has d-p<sub>z</sub> bonding characteristics, the Nb-As bonding in Nb<sub>2</sub>AsC exhibits obvious  $d-p_x + p_y$  bonding characteristics.

Based on the bonding features of Nb<sub>2</sub>AsC, its superior mechanical properties could be best understood in effectively enhanced Nb–As pd covalent bonding strength. Sufficient numbers of valence electrons provided by the As atoms cause the appearance of Nb d–As  $(p_x + p_y)$  bonding, and thereafter, the Nb d–As p bonding states shift downward to the same energy level as those of the Nb–C bond. The stronger the T–A covalent bonding in Nb<sub>2</sub>AsC, the higher resistance to applied strain is achieved. We have previously investigated mechanical properties of M<sub>2</sub>AlC (M = Ti, V, Cr, and Nb) by changing the VEC in compounds [3]. It was found that the M d–Al  $p_z$  bonding strengthened when VEC increased. However, after the sheardeformation resistive M d–Al  $p_z$  bonding states were completely occupied, filling metallic d–d bonding states on the M site occurred. In terms of mechanical parameters, shear moduli reach maximum values and drop slightly thereafter. Compared to the results in the above-mentioned references, tailoring the VEC of A-element atoms, rather than the T atoms in layered TAX carbides, may be more effective in optimizing the mechanical properties.

In summary, interesting bonding characteristics and superior mechanical properties of Nb<sub>2</sub>AsC are reported in the present letter. We found that Nb–As bonding states locate approximately in the same energy level as those of Nb–C bonding. This bonding characteristic is different from those of Al-containing ternary carbides, Ti<sub>2</sub>AlC, Nb<sub>2</sub>AlC, and Cr<sub>2</sub>AlC, wherein the T–Al and T–C (T = Ti, Nb and Cr) bonds extend in two subgroups and are separated by a pseudogap. The results suggest that Nb–As and Nb–C bonds have similar bonding strength, and thereafter the bulk modulus, shear modulus  $c_{44}$  and ideal shear strength of Nb<sub>2</sub>AsC are significantly enhanced. Finally, the present letter predicts underlying mechanisms toward strengthening ternary TAX carbides by tailoring the VEC of A-element atoms.

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